

Magnetostatics in Ponderable Media

Magnetization	$\vec{M} = \sum_i N_i \langle \vec{m}_i \rangle$
Magnetic permeability	μ
Magnetic susceptibility	$c_M = \partial M / \partial H$
Magnetic Field	$\vec{H} = \frac{1}{\mu_0} \vec{B} - \vec{M}$
Maxwell equations in media	$\vec{\nabla} \times \vec{H} = \vec{J} \quad \vec{\nabla} \cdot \vec{B} = 0$ isotropic & linear media $\vec{B} = \mu \vec{H} \quad \vec{M} = c_M \vec{H}$
Boundary conditions	$(\vec{B}_2 - \vec{B}_1) \cdot \vec{n}_{21} = 0$ $\vec{n}_{21} \times (\vec{H}_2 - \vec{H}_1) = \vec{a}$ or $\vec{B}_2 \cdot \vec{n}_{21} = \vec{B}_1 \cdot \vec{n}_{21} \quad \vec{B}_2 \times \vec{n}_{21} = \frac{\mu}{\mu_0} \vec{B}_1 \times \vec{n}_{21}$ $\vec{H}_2 \cdot \vec{n}_{21} = \frac{\mu}{\mu_0} \vec{H}_1 \cdot \vec{n}_{21} \quad \vec{H}_2 \times \vec{n}_{21} = \vec{H}_1 \times \vec{n}_{21}$
Magnetic Scalar Potential	$J = 0 \rightarrow \vec{H} = -\nabla \Phi_M$ linear media $\nabla^2 \Phi_M = -\mathbf{r}_M$ Effective magnetic charge density $\mathbf{r}_M = -\vec{\nabla} \cdot \vec{M}$
Solution for Scalar potential	$\Phi_M(\vec{r}) = -\frac{1}{4\pi} \int_V \frac{\vec{\nabla}' \cdot \vec{M}(\vec{r}')}{ \vec{r} - \vec{r}' } d^3r' + \frac{1}{4\pi} \oint_S \frac{\mathbf{n}' \cdot \vec{M}(\vec{r}')}{ \vec{r} - \vec{r}' } da'$
Energy of the field	$W = \frac{1}{2} \int \vec{H} \cdot \vec{B} d^3r = \frac{1}{2} \int \vec{J} \cdot \vec{A} d^3r$
Change in Energy	$W = \frac{1}{2} \int \vec{M} \cdot \vec{B}_0 d^3r$ putting object with $M \neq 0$ into magnetic field

Useful Relations

Spherical \leftrightarrow Cartesian

$$\hat{r} = \cos \theta \hat{x} + \sin \theta \hat{y} + \cos \theta \hat{z}$$

$$\hat{q} = \cos \theta \cos \phi \hat{x} + \sin \theta \cos \phi \hat{y} - \sin \theta \hat{z}$$

$$\hat{j} = -\sin \theta \hat{x} + \cos \theta \hat{y}$$

$$\hat{x} = \cos \theta \sin \phi \hat{r} + \cos \theta \cos \phi \hat{q} - \sin \theta \hat{j}$$

$$\hat{y} = \sin \theta \sin \phi \hat{r} + \sin \theta \cos \phi \hat{q} + \cos \theta \hat{j}$$

$$\hat{z} = \cos \theta \hat{r} - \sin \theta \hat{q}$$

Polar \leftrightarrow Cartesian

$$\hat{r} = \cos \theta \hat{x} + \sin \theta \hat{y} \quad \hat{q} = -\sin \theta \hat{x} + \cos \theta \hat{y}$$

$$\hat{x} = \cos \theta \hat{r} - \sin \theta \hat{q} \quad \hat{y} = \sin \theta \hat{r} + \cos \theta \hat{q}$$

Moving body current

$$\vec{J} = \mathbf{r} \cdot \vec{v}$$

Rotating body current

$$\vec{J} = \mathbf{r} \cdot (\vec{\omega} \times \vec{r})$$

Rotating hollow sphere

$$J_j = \omega R^2 \sin \theta \cdot \mathbf{d} \cdot (r - R)$$

Rotating hollow cylinder

$$J_j = \omega R^2 \cdot \mathbf{d} \cdot (r - R)$$

Infinite wire

$$B_j(r) = \frac{\mu_0 I}{2\pi r}$$

Force between 2 wires

$$\frac{dF}{dl} = \frac{\mu_0}{2\pi} \frac{I_1 I_2}{d}$$

Finite Solenoid

$$B_z = \frac{1}{2} \mu_0 N I (\cos \theta_1 + \cos \theta_2)$$

Time dependent fields

Magnetic flux	$F = \int_S \vec{B} \cdot \vec{n} da$
Electromotive force	$\mathbf{e} = \oint_C \vec{E} \cdot d\vec{l}$
Faraday induction law	$\mathbf{e} = -dF/dt$ $\vec{\nabla} \times \vec{E} + \frac{\partial \vec{B}}{\partial t} = 0$
Coefficients of Self- & Mutual Inductances	$W = \frac{1}{2} \sum_{i=1}^N L_i I_i^2 + \sum_{i=1}^N \sum_{j>i}^N M_{ij} I_i I_j$
Mutual inductance	$M_{ij} = F_{ij} / I_j = M_{ji}$ F_{ij} is flux from "j" linked with circuit "i"
EMF by inductance coefficients	$\mathbf{e}_{ij} = -dF_{ij}/dt$ Constant geometry $\mathbf{e}_{ij} = -M_{ij} \frac{dI_i}{dt} \quad \mathbf{e}_{ii} = -L_i \frac{dI_i}{dt}$ Constant currents $\mathbf{e}_{ij} = -I_j \frac{dM_{ij}}{dt} \quad \mathbf{e}_{ii} = -I_i \frac{dL_i}{dt}$

Maxwell Eq. and Gauge Transformations

	$\vec{\nabla} \cdot \vec{D} = \mathbf{r}$	$\vec{\nabla} \cdot \vec{B} = 0$
Maxwell Eq.	$\vec{\nabla} \times \vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t}$	$\vec{\nabla} \times \vec{E} + \frac{\partial \vec{B}}{\partial t} = 0$
	in vacuum, no sources	
	$\vec{\nabla} \cdot \vec{B} = 0$	$\vec{\nabla} \times \vec{B} = \mathbf{e}_0 \mu_0 \frac{\partial \vec{E}}{\partial t}$
	$\vec{\nabla} \cdot \vec{E} = 0$	$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$
Wave Equations	$\nabla^2 \vec{B} = \mu_0 \mathbf{e}_0 \frac{\partial^2 \vec{B}}{\partial t^2}$	$\nabla^2 \vec{E} = \mu_0 \mathbf{e}_0 \frac{\partial^2 \vec{E}}{\partial t^2}$
Velocity of light	$c = 1/\sqrt{\mu_0 \mathbf{e}_0}$	
Gauge Transformations	$\vec{A}' = \vec{A} + \nabla \Lambda$	$\Phi' = \Phi - \frac{\partial \Lambda}{\partial t}$
Delambertian	$\square^2 \equiv \nabla^2 - \frac{1}{c^2} \frac{\partial^2}{\partial t^2}$	
Lorenz gauge	$\vec{\nabla} \cdot \vec{A} + \frac{1}{c^2} \frac{\partial \Phi}{\partial t} = 0$	
	giving	$\square^2 \Phi = -\mathbf{r}/\mathbf{e}_0 \quad \square^2 \vec{A} = -\mu_0 \vec{J}$
Coulomb gauge	$\vec{\nabla} \cdot \vec{A} = 0$	
	giving	$\nabla^2 \Phi = -\mathbf{r}/\mathbf{e}_0 \quad \square^2 \vec{A} = -\mu_0 \vec{J}_t$
Poynting Vector	$\vec{S} = \vec{E} \times \vec{H}$	

Magnetostatics

Continuity equation	$\frac{\partial \mathbf{r}}{\partial t} + \vec{\nabla} \cdot \vec{J} = 0$
<i>in magnetostatics</i>	$\vec{\nabla} \cdot \vec{J} = 0$
Biot-Savart Law	$d\vec{B} = \frac{\mathbf{m}_0 I}{4\mathbf{p}} \frac{(d\vec{l} \times \vec{r})}{r^3}$
Moving Charge Field	$\vec{B} = \frac{\mathbf{m}_0 q}{4\mathbf{p}} \frac{\vec{v} \times \vec{r}}{r^3}$
First Ampere's Law	$d\vec{F} = I_1 (d\vec{l}_1 \times \vec{B})$
Force	$\vec{F} = \int \vec{J} \times \vec{B} d^3 r$
Torque	$\vec{N} = \int \vec{r} \times (\vec{J} \times \vec{B}) d^3 r$
Magnetic induction	$\vec{B}(\vec{r}) = \frac{\mathbf{m}_0}{4\mathbf{p}} \int \vec{J}(\vec{r}') \times \frac{\vec{r} - \vec{r}'}{ \vec{r} - \vec{r}' ^3} d^3 r'$
	$\vec{\nabla} \cdot \vec{B} = 0 \quad \vec{\nabla} \times \vec{B} = \mathbf{m}_0 \vec{J}$
Second Ampere's Law	$\oint_c \vec{B} \cdot d\vec{l} = \mathbf{m}_0 I$
Vector Potential	$\vec{A}(\vec{r}) = \frac{\mathbf{m}_0}{4\mathbf{p}} \int \frac{\vec{J}(\vec{r}')}{ \vec{r} - \vec{r}' } d^3 r' + \nabla \Psi(\vec{r})$
	$\vec{B} = \vec{\nabla} \times \vec{A} \quad \nabla(\vec{\nabla} \cdot \vec{A}) - \nabla^2 \vec{A} = \mathbf{m}_0 \vec{J}$
Magnetic moment	$\vec{m} = \frac{1}{2} \int \vec{r}' \times \vec{J}(\vec{r}') d^3 r'$
<i>current loop</i>	$\vec{m} = \frac{1}{2} \oint \vec{r} \times d\vec{l}$
<i>planar loop</i>	$m = I \cdot (\text{Area})$
<i>number of particles</i>	$\vec{m} = \frac{q}{2M} \vec{L}$
	where \vec{L} total angular momentum
Dipole Potential	$\vec{A}(\vec{r}) = \frac{\mathbf{m}_0}{4\mathbf{p}} \frac{\vec{m} \times \vec{r}}{r^3}$
Dipole induction	$\vec{B}(\vec{r}) = \frac{\mathbf{m}_0}{4\mathbf{p}} \left[\frac{3\vec{n}(\vec{n} \cdot \vec{m}) - \vec{m}}{r^3} \right]$
Force on dipole	$\vec{F} = \nabla(\vec{m} \cdot \vec{B})$
Torque on dipole	$\vec{N} = \int \left[(\vec{r}' \cdot \vec{B}) \vec{J} - (\vec{r}' \cdot \vec{J}) \vec{B} \right] d^3 r'$
Potential energy of dipole	$U = -\vec{m} \cdot \vec{B}$

Expanding Potential to Multipoles

Point charge expansion	$\frac{1}{ \vec{r} - \vec{r}' } = 4\mathbf{p} \sum_{l,m} \frac{1}{2l+1} \frac{r_{<}^l}{r_{>}^{l+1}} Y_{lm}^*(\mathbf{q}', \mathbf{f}) Y_{lm}(\mathbf{q}, \mathbf{f})$
Potential	$\Phi = \frac{1}{\mathbf{e}_0} \sum_{l,m} \frac{1}{2l+1} q_{lm} \frac{Y_{lm}(\mathbf{q}, \mathbf{f})}{r^{l+1}}$
or	$\Phi = \frac{1}{4\mathbf{p}\mathbf{e}_0} \left[\frac{q}{r} + \frac{\vec{p} \cdot \vec{r}}{r^3} + \frac{1}{2} \sum_{i,j} Q_{ij} \frac{x_i x_j}{r^5} + \dots \right]$
Multipoles	$q_{lm} = \int Y_{lm}^*(\mathbf{q}, \mathbf{f}) r^l \mathbf{r}(\vec{r}) d^3 r$
Dipole	$\vec{p} = \int \vec{r} \mathbf{r}(\vec{r}) d^3 r$
Dipole Field	$\vec{E} = \frac{3\vec{n}(\vec{p} \cdot \vec{n}) - \vec{p}}{4\mathbf{p}\mathbf{e}_0 \vec{r} - \vec{r}_0 ^3}$
Quadrupole	$Q_{ij} = \int (3x_i x_j - r^2 \mathbf{d}_{ij}) \mathbf{r}(\vec{x}) d^3 x$

* Multipole with order "n" is independent on the origin of coordinates only if all the multipoles with lower order vanish.

Electrostatics in Ponderable Media

Dipole moment density	$\vec{P} = N \langle \vec{p}_{mol} \rangle$
Electric Displacement	$\vec{D} = \mathbf{e}_0 \vec{E} + \vec{P}$
Electric susceptibility	$\mathbf{c}_e = \partial \mathbf{P} / \partial \mathbf{E}$
Dielectric constant	$\mathbf{e}/\mathbf{e}_0 = 1 + \mathbf{c}_e$
Electric permittivity	$\mathbf{e} = \mathbf{e}_0 (1 + \mathbf{c}_e)$
Maxwell equations in media	$\vec{\nabla} \times \vec{E} = 0 \quad \vec{\nabla} \cdot \vec{D} = \mathbf{r}$
<i>isotropic & linear media</i>	$\vec{D} = \mathbf{e} \vec{E} \quad \vec{P} = \mathbf{e}_0 \mathbf{c}_e \vec{E}$
<i>+ homogeneous media</i>	$\vec{\nabla} \cdot \vec{E} = \frac{\mathbf{r}}{\mathbf{e}}$
Boundary conditions	$(\vec{D}_2 - \vec{D}_1) \cdot \vec{n}_{21} = \mathbf{s} \quad E_{ 1} = E_{ 2}$
Induced charge density	$\mathbf{s}_{ind} = -(\vec{P}_2 - \vec{P}_1) \cdot \vec{n}_{21}$
Internal Field	$\vec{E}_i = \frac{\vec{P}}{3\mathbf{e}_0} + \vec{E}_{near}$
Average polarizability	$\langle \vec{p}_{mol} \rangle = \mathbf{e}_0 \mathbf{g}_{mol} (\vec{E} + \vec{E}_i)$
<i>Induced</i>	$\langle p_{mol} \rangle_{ind} = \frac{e^2}{m\nu_0^2} E$
<i>Permanent moment</i>	$\langle p_{mol} \rangle_0 \cong \frac{1}{3} \frac{p_0^2}{k_B T} E$
Molecular polarizability	$\mathbf{g}_{mol} = \frac{3 \mathbf{e}/\mathbf{e}_0 - 1}{N \mathbf{e}/\mathbf{e}_0 + 2}$
<i>Classical model</i>	$\mathbf{g}_{mol} \cong \mathbf{g}_{ind} + \frac{1}{3\mathbf{e}_0} \frac{p_0^2}{k_B T}$
Dielectrics Energy	$W = \frac{1}{2} \int \vec{E} \cdot \vec{D} d^3 r$
Change in Energy	$W = -\frac{1}{2} \int \vec{P} \cdot \vec{E}_0 d^3 r$
putting object with $M \neq 0$ into magnetic field	

Electrostatics

Electric field $\vec{E}(\vec{r}) = \frac{1}{4\pi\epsilon_0} \int \mathbf{r}(\vec{r}') \frac{\vec{r} - \vec{r}'}{|\vec{r} - \vec{r}'|^3} d^3r'$

Gauss Law $\int_S \vec{E} \cdot \vec{n} da = \frac{Q}{\epsilon_0} = \frac{1}{\epsilon_0} \int \mathbf{r}(\vec{r}) dV$
 $\vec{\nabla} \cdot \vec{E} = \frac{\rho}{\epsilon_0}$

Electric potential $\Phi(\vec{r}) = \frac{1}{4\pi\epsilon_0} \int \frac{\mathbf{r}(\vec{r}')}{|\vec{r} - \vec{r}'|} d^3r' + const$
 $E = -\nabla\Phi \quad \vec{\nabla} \times \vec{E} = 0$

Potential Energy

$$W = - \int_A^B \vec{F} \cdot d\vec{l} = q[\Phi(B) - \Phi(A)]$$

Electrostatic Potential Energy

$$W = \frac{1}{2} \int \mathbf{r}(\vec{r}) \Phi(\vec{r}) d^3r = \frac{\epsilon_0}{2} \int |\vec{E}|^2 d^3r$$

Energy density $w = \frac{\epsilon_0}{2} |\vec{E}|^2$

Electric force $\vec{F} = - \frac{\partial W}{\partial \vec{r}}$

Force per unit area in conductor $f = \frac{s^2}{2\epsilon_0}$

Boundary conditions $E_{\parallel 1} = E_{\parallel 2} \quad E_{\perp 2} - E_{\perp 1} = \frac{s}{\epsilon_0}$

Dipole layer potential $\Phi(\vec{r}) = - \frac{1}{4\pi\epsilon_0} \int_S D(\vec{r}') d\Omega'$

where $D(\vec{r}) = \lim_{d \rightarrow 0} \mathbf{s}(\vec{r}) \cdot d(\vec{r})$ & Ω is the solid angle from the observation point

Boundary: $\Phi_2 - \Phi_1 = \frac{D}{\epsilon_0}$

Poisson & Laplace Eq. $\nabla^2\Phi = -\rho/\epsilon_0 \quad \nabla^2\Phi = 0$

Point source $\nabla^2 \left(\frac{1}{|\vec{r} - \vec{r}'|} \right) = -4\pi\delta(\vec{r} - \vec{r}')$

Capacitance $C \equiv \frac{Q}{\Delta V}$

Orthogonal Function Proprieties

Legendre functions

$$P_l(x) = \frac{1}{2^l l!} \frac{d^l}{dx^l} (x^2 - 1)^l$$

$$\int_{-1}^1 P_l(x) P_l(x) dx = \frac{2}{2l+1} d_l \quad \int P_l dx = \frac{P_{l+1} - P_{l-1}}{2l+1}$$

$$P_l(1) = 1 \quad P_l(-1) = (-1)^l$$

$$P_l(0) = (-1)^{l/2} \frac{1 \cdot 3 \cdot 5 \cdots (l-1)}{2 \cdot 4 \cdot 6 \cdots n} \quad \text{n-even}$$

Spherical harmonics

$$Y_{lm}(\mathbf{p} - \mathbf{q}, \mathbf{f} + \mathbf{p}) = (-1)^l Y_{lm}(\mathbf{q}, \mathbf{f})$$

$$Y_{lm}(\mathbf{p} - \mathbf{q}, \mathbf{f}) = (-1)^{l+m} Y_{lm}(\mathbf{q}, \mathbf{f})$$

$$Y_{l,-m} = (-1)^m Y_{lm}^*$$

$$Y_{lm}(\mathbf{q}, \mathbf{f} + \mathbf{p}) = (-1)^m Y_{lm}(\mathbf{q}, \mathbf{f})$$

Green function

Green first identity

$$\int_V (f \nabla^2 y + \nabla f \cdot \nabla y) d^3x = \oint_S f \frac{\partial y}{\partial n} da$$

Green second identity (Green's Theorem)

$$\int_V (f \nabla^2 y - y \nabla^2 f) d^3x = \oint_S \left[f \frac{\partial y}{\partial n} - y \frac{\partial f}{\partial n} \right] da$$

Green Potential

$$\Phi(\vec{r}) = \frac{1}{4\pi\epsilon_0} \int_V \frac{\mathbf{r}(\vec{r}')}{|\vec{r} - \vec{r}'|} d^3r' + \frac{1}{4\pi} \oint_S \left[\frac{1}{|\vec{r} - \vec{r}'|} \frac{\partial \Phi}{\partial n'} - \Phi \frac{\partial}{\partial n'} \left(\frac{1}{|\vec{r} - \vec{r}'|} \right) \right]$$

Green Function

$$G(\vec{r}, \vec{r}') = \frac{1}{|\vec{r} - \vec{r}'|} + F(\vec{r}, \vec{r}') \quad \nabla'^2 F = 0$$

Dirichlet boundary $G_D = 0$ on S

$$\Phi(\vec{r}) = \frac{1}{4\pi\epsilon_0} \int_V \mathbf{r}(\vec{r}') G_D(\vec{r}, \vec{r}') d^3r' - \frac{1}{4\pi} \oint_S \Phi(\vec{r}') \frac{\partial G_D}{\partial n'} da'$$

Neumann boundary $\frac{\partial G_N}{\partial n'} = -\frac{4\pi}{S}$ on S

$$\Phi(\vec{r}) = \langle \Phi \rangle_S + \frac{1}{4\pi\epsilon_0} \int_V \mathbf{r}(\vec{r}') G_N(\vec{r}, \vec{r}') d^3r' + \frac{1}{4\pi} \oint_S G_N \frac{\partial \Phi}{\partial n'} da'$$

Green's reciprocity theorem

$$\int_V \mathbf{r} \Phi' d^3r + \int_S \mathbf{s} \Phi' da = \int_V \mathbf{r}' \Phi d^3r + \int_S \mathbf{s}' \Phi da$$

Orthogonal Functions Expansion

Rectangular $\Phi = e^{\pm ia_x x} e^{\pm ib_y y} e^{\pm z \sqrt{a^2 + b^2}}$

Polar $\Phi(r, f) = a_0 + b_0 \ln r + \sum_{n=1}^{\infty} (a_n r^n + b_n r^{-n}) (A \sin n f + B \cos n f)$

Spherical Symmetrical Azimuthally

$$\Phi(r, q) = \sum_{l=0}^{\infty} [A_l r^l + B_l r^{-(l+1)}] P_l(\cos q)$$

General Spherical

$$\Phi(r, q, f) = \sum_{l=0}^{\infty} \sum_{m=-l}^l [A_{lm} r^l + B_{lm} r^{-(l+1)}] Y_{lm}(q, f)$$

Point source $\frac{1}{|\vec{r} - \vec{r}'|} = 4\pi \sum_{l=0}^{\infty} \sum_{m=-l}^l \frac{1}{2l+1} \frac{r_{<}^l}{r_{>}^{l+1}} Y_{lm}^*(q', f') Y_{lm}(q, f)$

Low orders of Orthogonal functions

$$Y_{00} = \frac{1}{\sqrt{4\pi}} \quad Y_{20} = \sqrt{\frac{5}{4\pi}} \left(\frac{3}{2} \cos^2 q - \frac{1}{2} \right)$$

$$Y_{11} = -\sqrt{\frac{3}{8\pi}} \sin q e^{if} \quad Y_{21} = -\sqrt{\frac{15}{8\pi}} \sin q \cos q e^{if}$$

$$Y_{10} = \sqrt{\frac{3}{4\pi}} \cos q \quad Y_{22} = \frac{1}{4} \sqrt{\frac{15}{2\pi}} \sin^2 q e^{2if}$$

$$P_0 = 1 \quad P_1 = x$$

$$P_2 = \frac{1}{2} (3x^2 - 1) \quad P_3 = \frac{1}{2} (5x^3 - 3x)$$