

Electromagnetic Waves

Maxwell Relations $\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$ $\vec{\nabla} \times \vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t}$

$$\vec{\nabla} \cdot \vec{D} = 0 \quad \vec{\nabla} \cdot \vec{B} = 0$$

Electric Displacement $\vec{D} = \epsilon_0 \vec{E} + \vec{P} = \epsilon_r \epsilon_0 \vec{E}$

* true only in linear and uniform media

Magnetic Inductance $\vec{B} = \mu_0 (\vec{H} + \vec{M}) = \mu_r \mu_0 \vec{H}$

* true only in linear and uniform media

Velocity of Light $v = (\epsilon\mu)^{-\frac{1}{2}}$ $c = (\epsilon_0\mu_0)^{-\frac{1}{2}}$

Wave equations $\nabla^2 \vec{E} = \mu\epsilon \frac{\partial^2 \vec{E}}{\partial t^2}$ $\nabla^2 \vec{H} = \mu\epsilon \frac{\partial^2 \vec{H}}{\partial t^2}$

Wave number/frequency $k = \frac{2\pi}{\lambda}$ $\omega = kc = 2\pi f$

Index of Refraction $n = \frac{c}{v} = \sqrt{\epsilon}$

Plane Wave $\vec{E} = \vec{E}_0 \exp[i(\omega t - \vec{k} \cdot \vec{r})]$

Radial Wave $\vec{E} = \vec{E}_0 \frac{A}{r} \exp[i(\omega t - kr)]$

Medium Impedance $\frac{E}{H} = \sqrt{\frac{\mu}{\epsilon}} = Z$ $Z_0 = 377\Omega$

Field relations $\vec{k} \perp \vec{E}$ $\vec{k} \perp \vec{H}$ $\vec{H} \perp \vec{E}$

Poynting vector $\vec{S} = \vec{E} \times \vec{H}$ power per unit area

Poynting theorem flow of energy via closed surface

$$\oint_s (\vec{E} \times \vec{H}) \cdot d\vec{a} = \int_v \vec{E} \cdot \vec{J} + \frac{\partial}{\partial t} \left(\frac{\epsilon}{2} E^2 + \frac{\mu}{2} H^2 \right) + \vec{E} \cdot \frac{\partial \vec{P}}{\partial t} + \mu \vec{H} \cdot \frac{\partial \vec{M}}{\partial t} dV$$

Energy Consumption $\frac{1}{V} \langle P \rangle = \frac{1}{2} \omega \epsilon_0 |E|^2 \text{Im}(\chi_e)$

Wave Intensity $I = \langle \vec{S} \rangle$ if $T \gg \frac{2\pi}{\omega}$

while "T" is the interval between measurements

Boundary conditions $E_{1\parallel} = E_{2\parallel}$ $\epsilon_1 E_{1\perp} = \epsilon_2 E_{2\perp}$

$$H_{1\parallel} = H_{2\parallel} \quad \mu_1 H_{1\perp} = \mu_2 H_{2\perp}$$

Snell Law $n_i \sin \theta_i = n_t \sin \theta_t$

Transmittance/Reflection

$$r_{TE} = \frac{n_i \cos \theta_i - n_t \cos \theta_t}{n_i \cos \theta_i + n_t \cos \theta_t} \quad t_{TE} = \frac{2n_i \cos \theta_i}{n_i \cos \theta_i + n_t \cos \theta_t}$$

$$r_{TM} = \frac{n_t \cos \theta_i - n_i \cos \theta_t}{n_i \cos \theta_i + n_t \cos \theta_t} \quad t_{TM} = \frac{2n_i \cos \theta_i}{n_i \cos \theta_i + n_t \cos \theta_t}$$

$$T = t^2 \frac{n_t \cos \theta_t}{n_i \cos \theta_i} \quad R = r^2 \quad R + T = 1$$

Brewster Angle $\mathcal{G} = \arctan\left(\frac{n_t}{n_i}\right) \stackrel{n_t \approx 1.5}{=} 56^\circ$

Polarization Linear $E_{0x} = E_{0y}$ $\Delta\varphi = 0$

Circular $E_{0x} = E_{0y}$ $\Delta\varphi = \frac{\pi}{2}$

Elliptic else

Ray Optics

Conventions Distances from the right of the system are positive & from left are negative. Curvature radius with center from the right is positive. Angles beyond x-axis are positive.

v – distance from system to image

u – distance from system to object

Complex Matrices matrix of a system with 1,2...N components in a row is $M = M_N \dots M_2 \cdot M_1$

Determinant propriety $\det M = \begin{vmatrix} A & B \\ C & D \end{vmatrix} = \frac{n_{in}}{n_{out}}$

Transmission in media $\begin{pmatrix} r_{out} \\ \hat{r}_{out} \end{pmatrix} = \begin{pmatrix} 1 & x \\ 0 & 1 \end{pmatrix} \begin{pmatrix} r_{in} \\ \hat{r}_{in} \end{pmatrix}$

Planar Surface $\begin{pmatrix} r_{out} \\ \hat{r}_{out} \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & \frac{n_1}{n_2} \end{pmatrix} \begin{pmatrix} r_{in} \\ \hat{r}_{in} \end{pmatrix}$

Spherical Surface $\begin{pmatrix} r_{out} \\ \hat{r}_{out} \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ \left(\frac{n_1}{n_2} - 1\right) \frac{1}{R} & \frac{n_1}{n_2} \end{pmatrix} \begin{pmatrix} r_{in} \\ \hat{r}_{in} \end{pmatrix}$

Spherical Mirror $\begin{pmatrix} r_{out} \\ \hat{r}_{out} \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ \frac{2}{R} & 1 \end{pmatrix} \begin{pmatrix} r_{in} \\ \hat{r}_{in} \end{pmatrix}$

Thin lens $\begin{pmatrix} r_{out} \\ \hat{r}_{out} \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ -\frac{1}{f} & 1 \end{pmatrix} \begin{pmatrix} r_{in} \\ \hat{r}_{in} \end{pmatrix}$

Focal point $\frac{1}{f} = \frac{n_2 - n_1}{n_1} \left(\frac{1}{R_1} - \frac{1}{R_2} \right)$

Paths of Rays

- 1) ray through the center travels unchanged
- 2) parallel rays meet at the focal plane
- 3) parallel rays to the axis meet at the focus

Imaging if $\begin{pmatrix} r_{out} \\ \hat{r}_{out} \end{pmatrix} = \begin{pmatrix} A & B \\ C & D \end{pmatrix} \begin{pmatrix} r_{in} \\ \hat{r}_{in} \end{pmatrix}$ then $B = 0$

$$\begin{pmatrix} r_{out} \\ \hat{r}_{out} \end{pmatrix} = \begin{pmatrix} 1 & v \\ 0 & 1 \end{pmatrix} \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} 1 & u \\ 0 & 1 \end{pmatrix} \begin{pmatrix} r_{in} \\ \hat{r}_{in} \end{pmatrix}$$

Effective focus $\tilde{f} = -\frac{1}{c}$

Focal planes $F_2 = -\frac{a}{c} = a\tilde{f}$ $F_1 = -\frac{d}{c} = d\tilde{f}$

Principal planes $u_p = \frac{1-d}{c} = (d-1)\tilde{f}$ $v_p = \frac{1-a}{c} = (a-1)\tilde{f}$

planes for which the system acts like thin lens

Imaging condition $\frac{1}{u-u_p} + \frac{1}{v-v_p} = \frac{1}{f}$

Newton equation $\left(A - \frac{v}{f}\right) \left(D - \frac{u}{f}\right) = 1$

Magnification Linear $m = A = -\frac{v-v_p}{u-u_p}$
 $m = 1$ for $v = v_p$ $u = u_p$

Angle $D = \frac{1}{m}$

Image formation
 $v > 0$ Real & inverted image
 $v < 0$ Imaginary & straight image
 $f < 0$ Inverted image
 $f > 0$ Straight image

Approximation Conditions

Phase approximation $kd = k(d^2 + x^2 + y^2)^{1/2} \approx kd \left(1 + \frac{1}{2} \frac{r^2}{d^2}\right) + \dots$

d - distance to screen

r - distance off-axis of the aperture

Point source requirement $D\rho/d_1 < \frac{1}{4}\lambda$

D - effective diameter of the source

ρ - radius of the aperture

Paraxial approximation $\frac{x}{d}, \frac{y}{d} \ll 1$ $\sin x \approx x$

Maximum divergence angle $\alpha_{\max} = \arctan\left(\frac{a}{d}\right) \cong \frac{a}{d}$

“ a ” - radius of the diffraction pattern

Fresnel number $N_F \equiv \frac{a^2}{\lambda d}$

Inverse Fresnel number $N'_F \equiv \frac{b^2}{\lambda d}$

“ b ” - radius of the aperture

Fresnel diffraction condition $\frac{1}{4} N_F \alpha_{\max}^2 \ll 1$

Fraunhofer diffraction conditions $N'_F \ll 1$ $N_F \ll 1$

Interference

Two monochromatic waves $I = I_1 + I_2 + 2\sqrt{I_1 I_2} \cos \varphi$

Phase difference $\varphi = \mathbf{k}_1 \cdot \mathbf{r}_1 - \mathbf{k}_2 \cdot \mathbf{r}_2 + (\varepsilon_1 - \varepsilon_2)$

Equal Amplitude $I = 4I_0 \cos^2 \frac{\varphi}{2}$

Beating $I = I_1 + I_2 + 2\sqrt{I_1 I_2} \cos [2\pi(\nu_2 - \nu_1)t + \varepsilon(x, y)]$

Fringe velocity $\frac{dx}{dt} = -\frac{(d\varphi/dt)_x}{(d\varphi/dx)_t}$

Furrier Transform Proprieties

General proprieties

$$FT[f(x - x_0)] = e^{-2\pi i \nu x_0} F(\nu)$$

$$FT\left[f\left(\frac{x}{x_0}\right)\right] = |x_0| F(x_0 \nu)$$

$$FT[f(x)f(y)] = F(\nu_x)F(\nu_y)$$

$$FT[f_1 \otimes f_2] = F_1(\nu) \cdot F_2(\nu)$$

$$FT[FT[f(x, y)]] = f(-x, -y)$$

$f(x)$ symmetrical $\rightarrow F(-\nu) = F^*(\nu)$

$f(x)$ real & symmetrical $\rightarrow F(\nu)$ also

$$F(\bullet \bullet) = F(\bullet) \cdot F(\bullet \bullet)$$

Useful transforms

$$FT\left[rect\left(\frac{x}{a}\right)\right] = a \operatorname{sinc}(a\nu_x) \equiv a \frac{\sin(\pi a \nu_x)}{\pi a \nu_x}$$

while $rect\left(\frac{x}{a}\right) = 1$ $|x| \leq a/2$

$$FT\left[circ(r)\right] = \frac{J_1(2\pi \nu_p)}{\nu_p} \quad \text{where } \nu_p^2 = \nu_x^2 + \nu_y^2$$

Useful Fresnel Integrals on axis

$$\text{Circular hole} \quad I = |g|^2 = \frac{8A^2 \pi^2}{k^2} \left(1 - \cos \frac{k\rho^2}{2z_0}\right)$$

$$\text{Circular disc} \quad I = \frac{4A^2 \pi^2}{k^2}$$

Furrier Optics

Spatial Frequency $\nu_x = \frac{k_x}{2\pi} = \frac{x}{\lambda d}$ $\nu_y = \frac{k_y}{2\pi} = \frac{y}{\lambda d}$

Incident angles $\sin \theta_x = \lambda \nu_x$ $\sin \theta_y = \lambda \nu_y$

in paraxial approximation $\theta_x = \lambda \nu_x$ $\theta_y = \lambda \nu_y$

Spatial periods $\Lambda_x = \nu_x^{-1}$ $\Lambda_y = \nu_y^{-1}$ $\Lambda_z = \nu_z^{-1}$

Distraction by obstacle/Lens $\theta = \frac{\lambda}{d}$

Phase mask distraction if $f(x, y) = e^{-2\pi i \phi(x, y)}$

then $\nu_x(x) = \frac{\partial \phi}{\partial x}$ $\nu_y(y) = \frac{\partial \phi}{\partial y}$

Input-Output Relations in linear shift-invariant system (without magnification)

$$G(\nu_x, \nu_y) = H(\nu_x, \nu_y) F(\nu_x, \nu_y)$$

$$g(x, y) = f(x, y) \otimes h(x, y)$$

Transfer Function

$$\text{Free space} \quad H = \exp\left[-2\pi i \left(\frac{1}{\lambda^2} - \nu_x^2 - \nu_y^2\right)^{1/2} d\right]$$

$$\text{Far field} \quad \nu_p^2 \leq \lambda^{-2} \quad \rightarrow \quad |H| = 1$$

$$\text{where} \quad \nu_p^2 = \nu_x^2 + \nu_y^2$$

$$\text{Near field} \quad \nu_p^2 \geq \lambda^{-2} \quad |H| = e^{-2\pi d \sqrt{\frac{2}{\lambda}(\nu_p - \lambda^{-1})}}$$

here ν_p is a **Cut-Off** frequency

$$\text{Fresnel Approx.} \quad H = H_0 \exp\left[i\pi \lambda d (\nu_x^2 + \nu_y^2)\right]$$

where $H_0 = \exp(-ikd)$

Impulse-Response Function a response of the system to point source at the origin (δ func.).

Inverse FT of the Transfer function.

$$\text{Free Space} \quad h(x, y) = h_0 \exp\left[-ik \frac{x^2 + y^2}{2d}\right]$$

in Fresnel approx. while $h_0 = \frac{i}{\lambda d} e^{-ikd}$

Infinite Opening Lens

$$h(x, y) = h_1 h_2 \exp\left[-i \frac{k}{2f} (x^2 + y^2)\right] \delta\left(-\frac{x}{\lambda d_2}, -\frac{y}{\lambda d_2}\right)$$

$$\text{Finite Opening Lens} \quad h(x, y) = h_1 h_2 \hat{P}\left(-\frac{x}{\lambda d_2}, -\frac{y}{\lambda d_2}\right)$$

while \hat{P} is FT of the aperture (pupil) function and the varying phase was neglected.

Diffraction pattern width

$$\text{Circular:} \quad \Delta = 2r_1 = \frac{1.22 \lambda f}{D} \quad D - \text{mask diameter}$$

$$\text{Rectangular:} \quad \Delta_x = \frac{2\lambda f}{b_x} \quad b_x - \text{“x” mask width}$$

Furrier Transform by Lens

$$g(x, y) = \frac{i}{\lambda f} e^{-ik(f+d)} e^{i\pi \lambda (d-f) \frac{x^2 + y^2}{(\lambda f)^2}} F\left(\frac{x}{\lambda f}, \frac{y}{\lambda f}\right)$$

plane wave is focused at $x_0 = \lambda f \nu_x$ $y_0 = \lambda f \nu_y$

Fraunhofer Diffraction

$$g(x, y) = h_0 \exp\left[-i \frac{\pi}{\lambda d} (x^2 + y^2)\right] F\left(\frac{x}{\lambda d}, \frac{y}{\lambda d}\right)$$